

Dinamica dell'Electron Cloud: Calcolo dei Coefficienti della Mappa

T. Demma¹, S.Petracca^{2,3}, **A. Stabile²**

¹ LNF – INFN, Frascati

² Dipartimento di Ingegneria - Universita' del Sannio - Benevento

³ Sezione INFN di Napoli, gruppo collegato di Salerno

SOCIETÀ ITALIANA DI FISICA
XCVI° CONGRESSO NAZIONALE
Bologna, 20 Settembre 2010



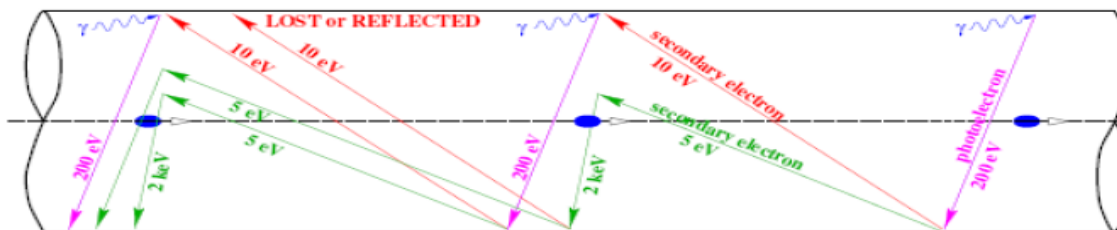


Plan of Talk

- Electron Cloud Effect
- Cubic Map Formalism
- Saturation Condition and the Quadratic Coefficient
- Conclusions and Outlook



Electron Cloud Effect and Secondary Emission Yield



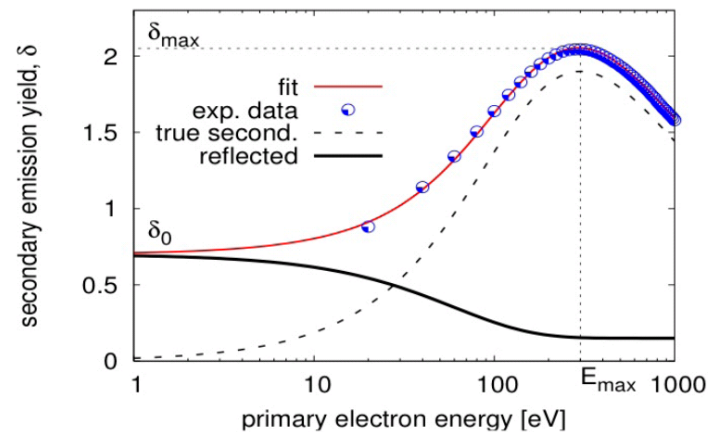
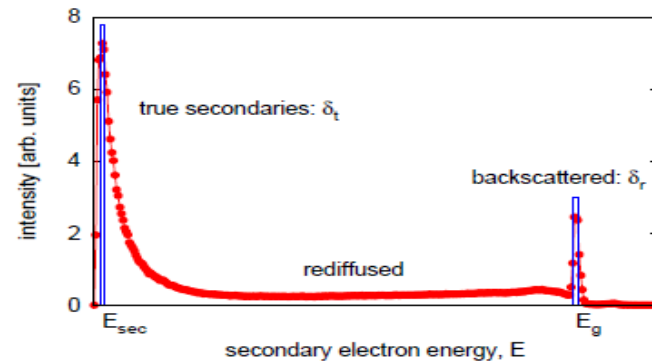
The electron cloud develops quickly as photons striking the vacuum chamberwall knock out electrons that are then accelerated by the beam, gain energy, and strike the chamber again, producing more electrons.

The interaction between the electron cloud and a beam leads to the electron cloud effects such as single- and multi-bunch instability, tune shift, increase of pressure and so on.

$$\delta = I_s/I_p \text{ and } \delta = \delta_{ts} + \delta_{el} + \delta_{rdf}$$

$$\delta = \delta_{ts} + \delta_{el}$$

Typical vacuum material have $\delta_{max} > 1$

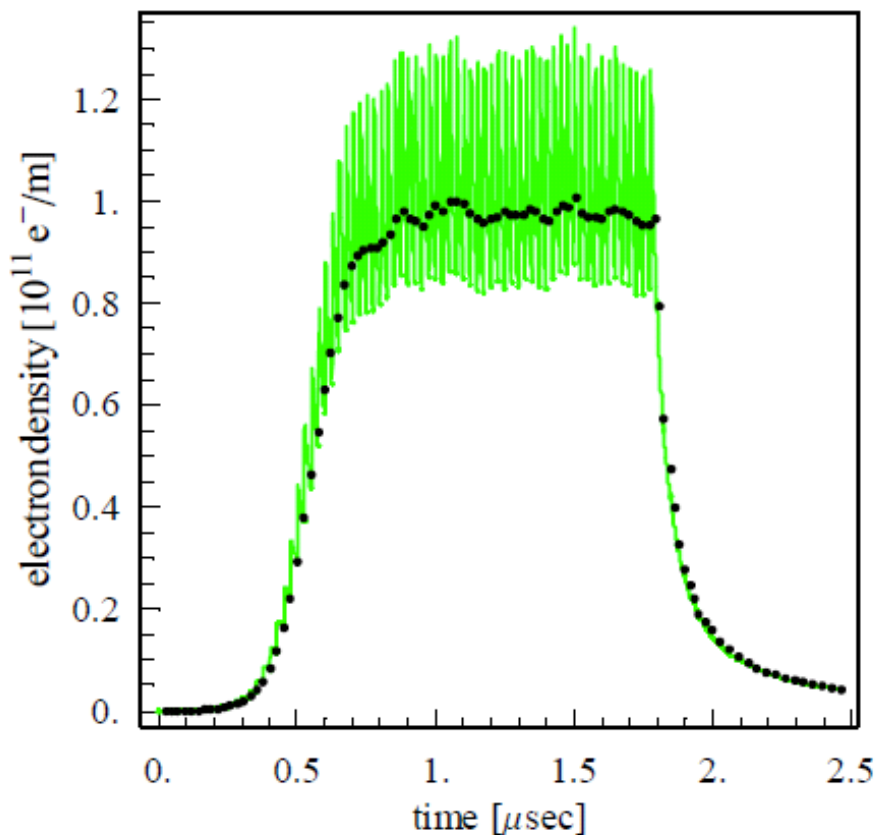




Electron Cloud Buildup I

$$\delta_{max} = 1.7 \text{ (72 f, 28 e)}$$

LHC dipole parameters



parameter	unit	value
beam particle energy	GeV	7000
bunch spacing	ns	25
bunch length	m	0.075
number of bunches N_b	-	72
number of particles per bunch N	10^{11}	0.8 to 1.6
bending field B	T	8.4
length of bending magnet	m	14.2
vacuum screen half height	m	0.018
vacuum screen half width	m	0.022
circumference	m	27000
primary photo-emission yield	-	$7.98 \cdot 10^{-4}$
maximum SEY δ_{max}	-	1.3 to 1.7
energy for max. SEY E_{max}	eV	237.125
energy width for secondary e^-	eV	1.8

— ECLLOUD (CERN) output

- bunch to bunch average



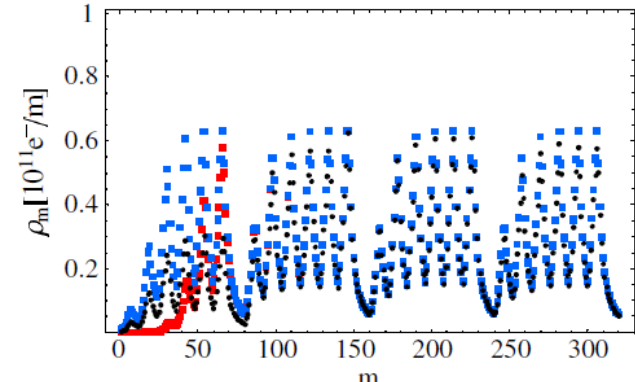
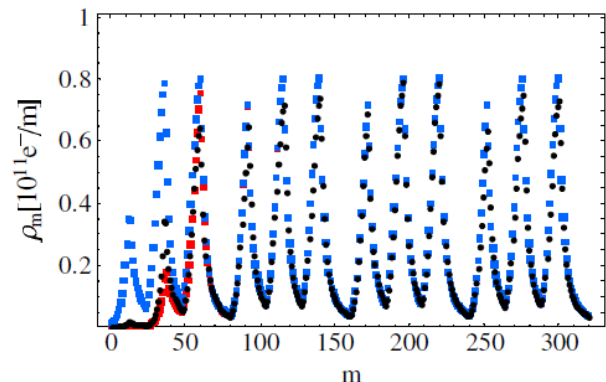
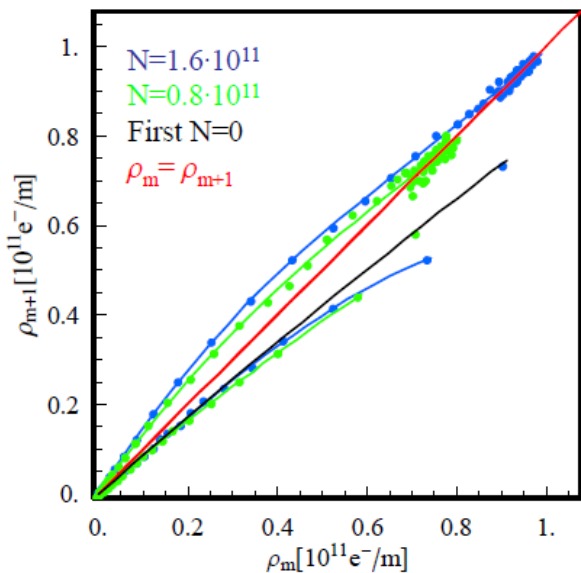
Electron Cloud Buildup II

Bunch Filling Patterns

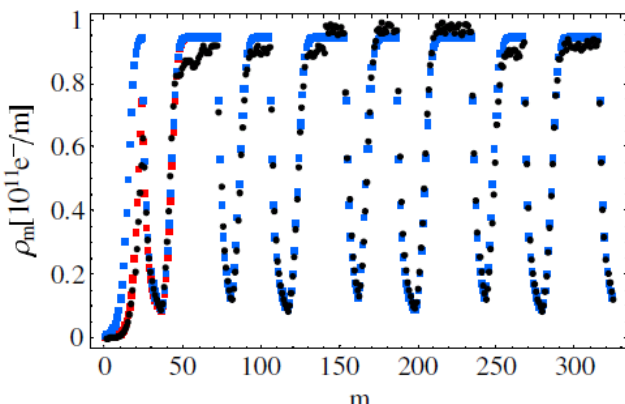
$\delta_{max} = 1.7$ (72 f, 28 e)

(12f,12e, 12f,12e,12f,12e)

(6f,6e)



(24f, 12e, 36 f)



$\rho_0 = 10^{-2} [10^{11} e^-/m]$, $\rho_0 = 10^{-4} [10^{11} e^-/m]$

Map results do not depends on the initial electron density

Relative error below 20%

Different bunch patterns can be described using the same Map coefficients

Lines corresponds to cubic fit of the form:

$$n_{m+1} = \alpha n_m + \beta n_m^2 + \gamma n_m^3$$

Three sets of coefficients are needed to describe the e-cloud density evolution



Maps and Dynamics for Electron Clouds

- For a given beam pipe characteristics (SEY, chamber dimensions, etc.) the evolution of the electron density is only driven by the bunch passing by, and the existing electron density before the bunch passage:

$$n_{m+1} = F(n_m)$$

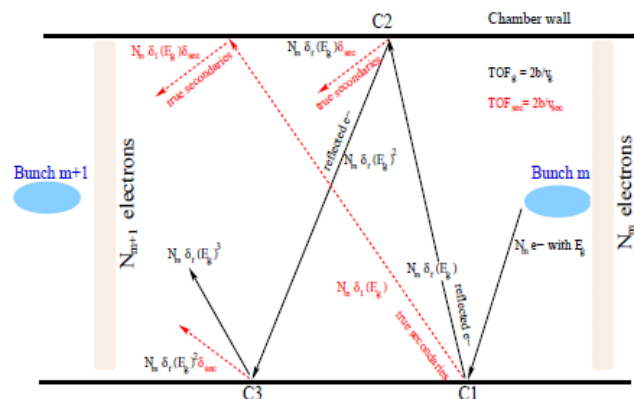
- Simplify the e-cloud problem into a small number of mathematical parameters.
- The bunch-to-bunch evolution of the e-cloud density is represented by a cubic map:

$$n_{m+1} = \alpha n_m + \beta n_m^2 + \gamma n_m^3$$

where n_m is the bunch to bunch average of the electron line density.

- The linear term a describes the exponential growth of the cloud density. For $n \gg 1$ $n_1 = \alpha n_0$, $n_2 = \alpha n_1 = \alpha^2 n_0$, and $n_m = \alpha n_{m-1} = \alpha^m n_0$
- The quadratic term b is related to the saturation due to space charge
- The cubic term is always very small and could be interpreted additional correction embodying small corrections.

$$a = \frac{N_{m+1}}{N_m} = \delta_r^n(\bar{E}_g) + \delta_t(\bar{E}_g)\delta_s^\eta(E_s) \cdot \frac{\delta_s^\eta(E_s) - \delta_r^n(\bar{E}_g)}{\delta_s^\eta(E_s) - \delta_r(\bar{E}_g)}$$





Saturation Condition

In free field regions the distance (in units of beampipe radius b) passed by secondary electrons (with energy ε_0) before the next bunch arrives is

$$\xi = \frac{h}{b} \sqrt{\frac{2\varepsilon_0}{mc^2}}$$

The density of the secondary electrons grows until the space-charge potential is lower than energy of the secondary electrons. For a gausslike distribution of electron cloud

$$V(r) = \int_r^b \vec{E} \cdot d\vec{l} = -\frac{\bar{N}_b e}{2\pi\epsilon_0 h} \ln \frac{r}{b} - \frac{N_{el} e}{2\pi\epsilon_0 h} \frac{1}{\int_a^b e^{-\frac{(y-r_0)^2}{2\sigma^2}} y dy} \int_a^y \frac{e^{-\frac{(z-r_0)^2}{2\sigma^2}} z dz}{y} dy = \quad g = \bar{N}_b / N_{el}$$

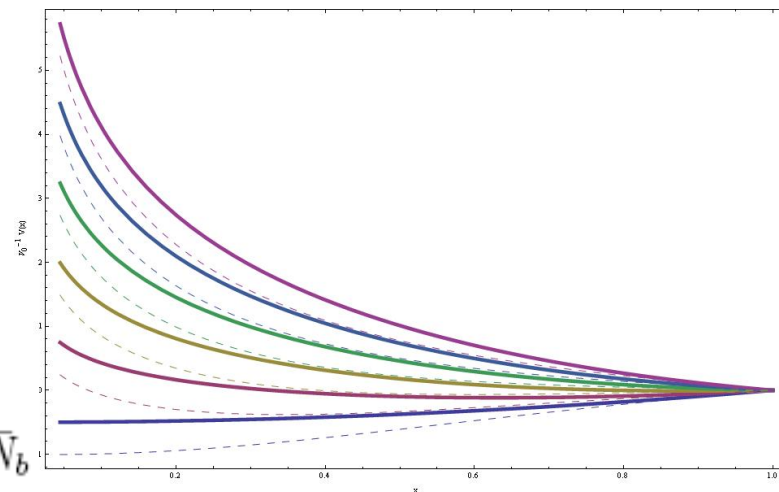
$$-\frac{\bar{N}_b e}{2\pi\epsilon_0 h} \ln x - \frac{N_{el} e}{2\pi\epsilon_0 h} \frac{G(x)}{F(1)} = -V_0 \left[g \ln x + \frac{G(x)}{F(1)} \right]$$

If the distribution is uniform

$$V_{ud}(r) = -V_0 \left[g \ln x + \frac{1-x^2}{2} \right]$$

The saturation condition can be obtained requiring

$$-eV(1-\xi) \sim \varepsilon_0 \quad N_{el,sat} = \frac{2\pi\epsilon_0 h F(1) \varepsilon_0}{e^2 G(1-\xi)} - \frac{F(1) \ln(1-\xi)}{G(1-\xi)} \bar{N}_b$$





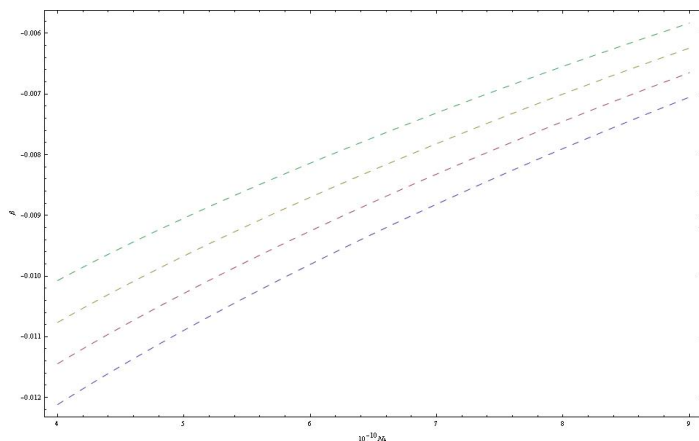
Electron density and quadratic coefficient

- The critical density is given by relation
- The saturation condition satisfies the equation
- The quadratic map becomes

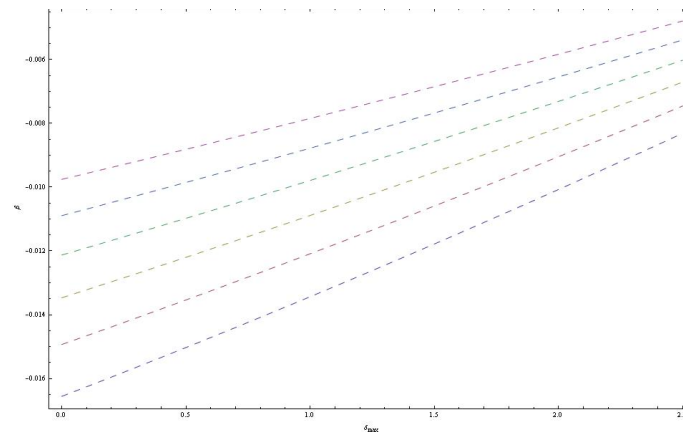
$$n_{sat} = \frac{N_{el,sat}}{\pi h b^2 [(\tilde{r}_0 + p \tilde{\sigma})^2 - \tilde{a}^2]}$$

$$n_{sat} = \alpha n_{sat} + \beta n_{sat}^2 \rightarrow \beta = \frac{1 - \alpha}{n_{sat}}$$

$$n_{m+1} = \alpha n_m + \frac{1 - \alpha}{n_{sat}} n_m^2$$



Analytical prediction of coefficient for values $\delta_{max} = 1,4 - 2$ and $p = 2$.



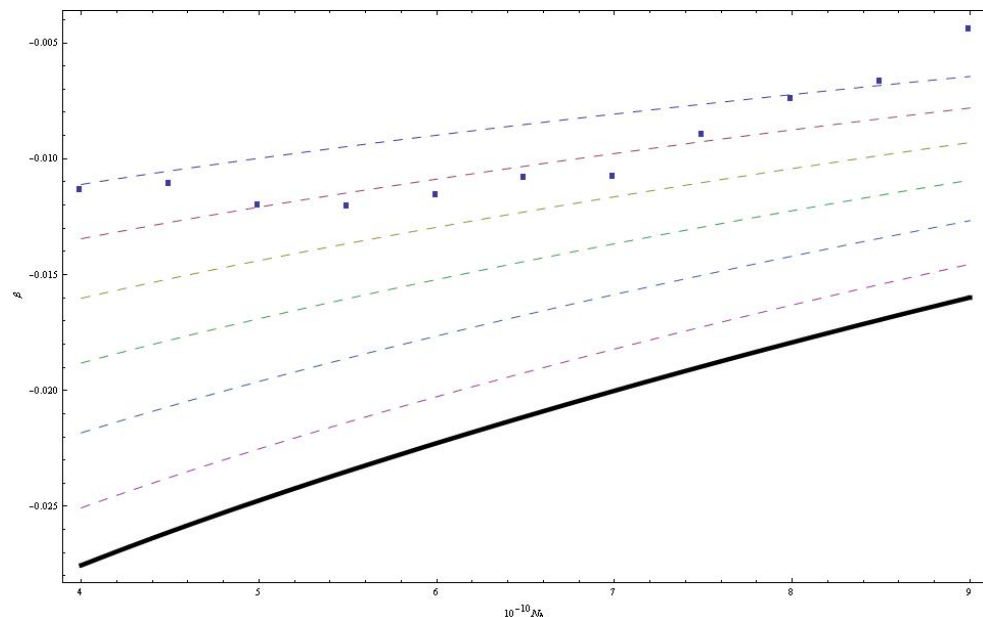
Analytical prediction of coefficient for values $N_b = (4 - 9) * 10^{10}$ and $p = 2$.



Quadratic coefficient

Comparison of the quadratic coefficient derived using ECLLOUD simulations (points) and using the analysis of previous slides (dashed lines) with $p = 2 - 3$. The solid line is the result by assuming an uniform density.

Parameter	Quantity	Unit	Value
Beam pipe radius	b	m	.045
Beam size	a	m	.002
Bunch spacing	s_b	m	1.2
Bunch length	h	m	.013
Energy for δ_{max}	$\mathcal{E}_{0,max}$	eV	300
-	\mathcal{E}_r	eV	60
Particles per bunch	N_b	10^{10}	$4 \div 9$
SEY (max)	δ_{max}	-	1.7
SEY ($\mathcal{E} \rightarrow 0$)	δ_0	-	.7
SEY ($\mathcal{E} \rightarrow \infty$)	δ_∞	-	.15
-	ζ	-	1.83



Conclusions and Outlook

- The electron-cloud buildup can be described by a cubic map.
- Remarkably, if all other parameters (namely, the bunch charge N , the SEY, and the pipe parameters) are held fixed, the map coefficients basically *do not* depend on the filling pattern.
- An approximate formula has been derived for the quadratic coefficient in the map. The results are in acceptable agreement with numerical simulations obtained from E-CLOUD.
- The analytical result could be useful to determine safe regions in parameters space where to minimize the electron clouds.
- Furthermore we would extend our results for the quadratic coefficient in order to include the presence of a magnetic field.
- Work is in order to calculate the higher order terms in the map.

References



1. T. Demma, S. Petracca, F. Ruggiero, G. Rumolo, F. Zimmermann - Phys. Rev. ST Accel. Beams **10**, 14401 – 2007
2. U. Iriso, S. Pegg - Phys. Rev. ST Accel. Beams **8**, 024403 – 2005
3. U. Iriso, S. Pegg - Proceedings of EPAC – 2006
4. T. Demma, S. Petracca - Proceedings of EPAC 2008
5. T. Demma, R. Cimino, A. Drago, S. Petracca, AS - Proceedings of PAC 2009
6. T. Demma, S. Petracca, AS - Proceedings of IPAC 2010
7. T. Demma, S. Petracca, AS - Proceedings of RiNEm 2010